

# New Integral Representations for the Bessel Functions

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## Abstract

In this paper, new integral representations for the Bessel  $J$  and  $I$  functions were presented and their results were used to derive an expression for the Modified Bessel  $K$  function.

**Keywords:** Bessel functions, Modified Bessel functions, Hypergeometric function, Partial differential equation, Gamma function, Legendre's duplication formula, beta functions.

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## 1 Introduction

The theory of Bessel functions is intimately connected with the theory of a certain type of differential equation of the first order, known as Riccati's equation. In fact, a Bessel function is usually defined as a particular solution of a linear differential equation of the second order (known as Bessel's equation) which is derived from Riccati's equation by elementary transformation. [1]

Bessel functions, first defined by the mathematician Daniel Bernoulli and then generalized by Friedrich Bessel, are canonical solutions  $y(x)$  of Bessel's differential equation

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - \alpha^2)y = 0$$

Bessel functions of the first kind, denoted as  $J_\alpha(x)$ , are solutions of Bessel's differential equation. For integer or positive  $\alpha$ , Bessel functions of the first kind are finite at the origin; while for negative non-integer  $\alpha$ , Bessel functions of the first kind diverge as  $x$  approaches zero. It is possible to define the function by its series expansion around  $x = 0$ , which can be found by applying the Frobenius method to Bessel's equation

$$J_\alpha(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m + \alpha} \quad \forall x \in \mathbb{C}$$

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where  $\Gamma(z)$  is the gamma function, a shifted generalization of the factorial function to non-integer values. [5]

## 2 Integral representations for $J_\alpha(x)$ and $J_0(x)$

$$\begin{aligned}
 J_\alpha(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m+\alpha} \\
 J_\alpha(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m+\alpha} \\
 &= \frac{1}{\Gamma(\alpha)} \sum_{m=0}^{\infty} \frac{(-1)^m \Gamma(m+1) \Gamma(\alpha)}{m! \Gamma(m+1) \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m+\alpha} \\
 &= \frac{1}{\Gamma(\alpha)} \sum_{m=0}^{\infty} \frac{(-1)^m \Gamma(m+1) \Gamma(\alpha)}{(m!)^2 \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m+\alpha} \\
 &= \frac{1}{\Gamma(\alpha)} \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} \beta(m+1, \alpha) \left(\frac{x}{2}\right)^{2m+\alpha} \\
 &= \frac{1}{\Gamma(\alpha)} \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} \int_0^1 t^m (1-t)^{\alpha-1} dt \left(\frac{x}{2}\right)^{2m+\alpha} \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} t^m \left(\frac{x}{2}\right)^{2m} (1-t)^{\alpha-1} dt \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} \left(\frac{x\sqrt{t}}{2}\right)^{2m} (1-t)^{\alpha-1} dt
 \end{aligned}$$

We notice that

$$\begin{aligned}
 J_0(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + 0 + 1)} \left(\frac{x}{2}\right)^{2m+0} \\
 &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + 1)} \left(\frac{x}{2}\right)^{2m} \\
 &= \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} \left(\frac{x}{2}\right)^{2m} \\
 \implies J_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 J_0(x\sqrt{t}) (1-t)^{\alpha-1} dt \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 J_0(x\sqrt{1-t}) t^{\alpha-1} dt \tag{1}
 \end{aligned}$$

By Legendre’s duplication formula,

$$\begin{aligned}
 \Gamma\left(m + \frac{1}{2}\right) &= \frac{\Gamma(2m)\sqrt{\pi}}{\Gamma(m)2^{2m-1}} \\
 \implies \frac{\Gamma\left(m + \frac{1}{2}\right)2^{2m-1}}{\Gamma(2m)\sqrt{\pi}} &= \frac{1}{\Gamma(m)} \\
 \implies J_0(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + 1)} \left(\frac{x}{2}\right)^{2m}
 \end{aligned}$$

$$\begin{aligned}
 &= \sum_{m=0}^{\infty} \frac{(-1)^m}{\Gamma(m+1)^2} \left(\frac{x}{2}\right)^{2m} \\
 &= \sum_{m=1}^{\infty} \frac{(-1)^m}{\Gamma(m)\Gamma(m)} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} \frac{(-1)^{m-1}\Gamma(m+\frac{1}{2})2^{2m-1}}{\Gamma(2m)\Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} \frac{2m(-1)^{m-1}\Gamma(m+\frac{1}{2})2^{2m-1}}{2m\Gamma(2m)\Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} = \sum_{m=1}^{\infty} \frac{2m(-1)^{m-1}\Gamma(m+\frac{1}{2})2^{2m-1}}{\Gamma(2m+1)\Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} \frac{2m(-1)^{m-1}\Gamma(m+\frac{1}{2})\Gamma(m+\frac{1}{2})}{\Gamma(2m+1)} \cdot \frac{2^{2m-1}}{\Gamma(m+\frac{1}{2})\Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} 2m(-1)^{m-1}\beta\left(m+\frac{1}{2}, m+\frac{1}{2}\right) \cdot \frac{2^{2m-1}}{\Gamma(m+\frac{1}{2})\Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} 2m(-1)^{m-1}\beta\left(m+\frac{1}{2}, m+\frac{1}{2}\right) \cdot \frac{2^{2m-1}}{\Gamma(2m)\sqrt{\pi} \cdot \Gamma(m)\sqrt{\pi}} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \sum_{m=1}^{\infty} 2m(-1)^{m-1}\beta\left(m+\frac{1}{2}, m+\frac{1}{2}\right) \cdot \frac{2^{4m-2}}{\pi\Gamma(2m)} \left(\frac{x}{2}\right)^{2m-2} \\
 &= \frac{1}{2\pi} \sum_{m=1}^{\infty} m(-1)^{m-1}\beta\left(m+\frac{1}{2}, m+\frac{1}{2}\right) \cdot \frac{1}{\Gamma(2m)} (2x)^{2m} \left(\frac{x}{2}\right)^{-2} \\
 &= \frac{1}{2\pi} \sum_{m=1}^{\infty} m(-1)^{m-1} \int_0^1 \omega^{m-\frac{1}{2}} (1-\omega)^{m-\frac{1}{2}} d\omega \cdot \frac{1}{\Gamma(2m)} (2x)^{2m} \left(\frac{x}{2}\right)^{-2} \\
 &= \frac{2}{\pi x^2} \int_0^1 \left( \omega^{-\frac{1}{2}} (1-\omega)^{-\frac{1}{2}} \sum_{m=1}^{\infty} m(-1)^{m-1} \omega^m (1-\omega)^m \cdot \frac{1}{\Gamma(2m)} (2x)^{2m} \right) d\omega \\
 &= \frac{2}{\pi x^2} \int_0^1 \left( \omega^{-\frac{1}{2}} (1-\omega)^{-\frac{1}{2}} \sum_{m=1}^{\infty} \frac{(-1)^{m-1} m (4x^2 (\omega - \omega^2))^m}{\Gamma(2m)} \right) d\omega
 \end{aligned}$$

Consider

$$\begin{aligned}
 &\sum_{m=1}^{\infty} \frac{(-1)^{m-1} m z^m}{\Gamma(2m)} \\
 &= z \sum_{m=1}^{\infty} \frac{(-1)^{m-1} \frac{\partial}{\partial z} (z^m)}{\Gamma(2m)} = z \frac{\partial}{\partial z} \left( \sum_{m=1}^{\infty} \frac{(-1)^{m-1} z^m}{\Gamma(2m)} \right) = z \frac{\partial}{\partial z} \left( \sum_{m=0}^{\infty} \frac{(-1)^m z^{m+1}}{\Gamma(2m+2)} \right) \\
 &= z \frac{\partial}{\partial z} \left( \sum_{m=0}^{\infty} \frac{(-1)^m z^{m+1}}{(2m+1)!} \right) = z \frac{\partial}{\partial z} \left( \sqrt{z} \sum_{m=0}^{\infty} \frac{(-1)^m (\sqrt{z})^{2m+1}}{(2m+1)!} \right) \\
 &= z \frac{\partial}{\partial z} (\sqrt{z} \sin \sqrt{z}) = z \left( \frac{1}{2\sqrt{z}} \sin \sqrt{z} + \frac{1}{2} \cos \sqrt{z} \right) = \frac{1}{2} (\sqrt{z} \sin \sqrt{z} + z \cos \sqrt{z})
 \end{aligned}$$

$$\begin{aligned}
 \Rightarrow J_0(x) &= \frac{2}{\pi x^2} \int_0^1 \left( \omega^{-\frac{1}{2}} (1-\omega)^{-\frac{1}{2}} \sum_{m=1}^{\infty} \frac{(-1)^m m (4x^2 (\omega - \omega^2))^m}{\Gamma(2m)} \right) d\omega \\
 &= \frac{2}{\pi x^2} \cdot \frac{1}{2} \int_0^1 \left( \omega^{-\frac{1}{2}} (1-\omega)^{-\frac{1}{2}} \left( \sqrt{4x^2 \omega(1-\omega)} \sin \left( \sqrt{4x^2 \omega(1-\omega)} \right) \right. \right. \\
 &\quad \left. \left. + 4x^2 \omega(1-\omega) \cos \left( \sqrt{4x^2 \omega(1-\omega)} \right) \right) \right) d\omega \\
 &= \frac{1}{\pi x^2} \int_0^1 2x \sin \left( \sqrt{4x^2 \omega(1-\omega)} \right) + 4x^2 \sqrt{\omega(1-\omega)} \cos \left( \sqrt{4x^2 \omega(1-\omega)} \right) d\omega
 \end{aligned}$$

On the transformation  $\omega \rightarrow \sin^2 \theta$ ,

$$\begin{aligned}
 J_0(x) &= \frac{1}{\pi x^2} \int_0^{\frac{\pi}{2}} \left( 2x \sin(2x \sin \theta \cos \theta) + 4x^2 \sin \theta \cos \theta \cos(2x \sin \theta \cos \theta) \right) 2 \sin \theta \cos \theta d\theta \\
 \implies J_0(x) &= \frac{1}{\pi x^2} \int_0^{\frac{\pi}{2}} \left( 2x \sin(2x \sin \theta \cos \theta) + 4x^2 \sin \theta \cos \theta \cos(2x \sin \theta \cos \theta) \right) 2 \sin \theta \cos \theta d\theta \\
 &= \frac{1}{\pi x^2} \int_0^{\frac{\pi}{2}} \left( 2x \sin(x \sin 2\theta) + 2x^2 \sin 2\theta \cos(x \sin 2\theta) \right) \sin 2\theta d\theta \\
 &= \frac{1}{\pi x^2} \int_0^\pi \left( x \sin(x \sin \phi) + x^2 \sin \phi \cos(x \sin \phi) \right) \sin \phi d\phi \\
 &= \frac{1}{\pi x} \int_0^\pi \left( \sin(x \sin \phi) + x \sin \phi \cos(x \sin \phi) \right) \sin \phi d\phi
 \end{aligned}$$

Therefore,

$$\begin{aligned}
 J_0(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} & (2) \\
 &= \frac{1}{\pi x} \int_0^\pi \left( \sin(x \sin \phi) + x \sin \phi \cos(x \sin \phi) \right) \sin \phi d\phi & (3)
 \end{aligned}$$

$$\begin{aligned}
 J_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 J_0(x\sqrt{1-t}) t^{\alpha-1} dt \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} 2 \sin^{2\alpha-1} \vartheta \cos \vartheta J_0(x \cos \vartheta) d\vartheta \\
 &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \frac{1}{\cos \vartheta} \left[ 2 \sin^{2\alpha-1} \vartheta \sin \phi \cos \vartheta [\sin(x \cos \vartheta \sin \phi) \right. \\
 &\qquad \qquad \qquad \left. + x \cos \vartheta \sin \phi \cos(x \cos \vartheta \sin \phi) \right] d\phi d\vartheta \\
 &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \left( 2 \sin^{2\alpha-1} \vartheta \sin \phi [\sin(x \cos \vartheta \sin \phi) + x \cos \vartheta \sin \phi \cos(x \cos \vartheta \sin \phi)] \right) d\phi d\vartheta
 \end{aligned}$$

### 3 Using the new integral representation to establish Lipschitz's result

It was shown by Lipschitz that [2]

$$\int_0^\infty e^{-at} J_0(bt) dt = \frac{1}{\sqrt{a^2 + b^2}}$$

The simplest method of establishing this result is to replace the Bessel coefficient by Parseval's integral representation and then change the order of integration by Fubini's theorem.

However, implementing the integral representation I discovered, Lipschitz's result can also be established.

$$\begin{aligned}
 J_0(x) &= \frac{1}{\pi x} \int_0^\pi \left( \sin(x \sin \phi) + x \sin \phi \cos(x \sin \phi) \right) \sin \phi d\phi \\
 \implies J_0(bt) &= \frac{1}{\pi bt} \int_0^\pi \left( \sin(bt \sin \phi) + bt \sin \phi \cos(bt \sin \phi) \right) \sin \phi d\phi
 \end{aligned}$$

$$\begin{aligned} \int_0^\infty e^{-at} J_0(bt) dt &= \int_0^\infty \frac{e^{-at}}{\pi bt} \int_0^\pi (\sin(bt \sin \phi) + bt \sin \phi \cos(bt \sin \phi)) \sin \phi d\phi dt \\ &= \int_0^\pi \int_0^\infty e^{-at} \left( \frac{\sin(bt \sin \phi)}{\pi bt} + \frac{1}{\pi} \sin \phi \cos(bt \sin \phi) \right) \sin \phi dt d\phi \\ &= \int_0^\pi \left( \int_0^\infty e^{-at} \frac{\sin(bt \sin \phi) \sin \phi}{\pi bt} dt + \int_0^\infty \frac{e^{-at}}{\pi} \sin^2 \phi \cos(bt \sin \phi) dt \right) d\phi \end{aligned}$$

Let

$$\Phi(b, \phi) = \int_0^\infty e^{-at} \frac{\sin(bt \sin \phi)}{\pi bt} dt$$

$$\begin{aligned} \Phi(b, \phi) &= \frac{1}{\pi b} \int_0^\infty e^{-at} \frac{\sin(bt \sin \phi)}{t} dt \\ &= \frac{1}{\pi b} \int_0^\infty e^{-at} \int_0^{b \sin \phi} \cos(\tau t) d\tau dt \\ &= \frac{1}{\pi b} \int_0^{b \sin \phi} \int_0^\infty e^{-at} \cos(\tau t) dt d\tau \\ &= \frac{1}{\pi b} \int_0^{b \sin \phi} \mathcal{L}\{\cos(\tau t)\}(a) d\tau \\ &= \frac{1}{\pi b} \int_0^{b \sin \phi} \frac{a}{\tau^2 + a^2} d\tau \\ &= \frac{1}{\pi b} \arctan\left(\frac{\tau}{a}\right) \Big|_0^{b \sin \phi} \\ &= \frac{1}{\pi b} \arctan\left(\frac{b \sin \phi}{a}\right) \end{aligned}$$

Let

$$\eta(b, \phi) = \int_0^\infty \frac{e^{-at}}{\pi} \cos(bt \sin \phi) dt$$

$$\begin{aligned} \eta(b, \phi) &= \int_0^\infty \frac{e^{-at}}{\pi} \cos(bt \sin \phi) dt \\ &= \frac{1}{\pi} \int_0^\infty e^{-at} \cos(bt \sin \phi) dt \\ &= \frac{1}{\pi} \mathcal{L}\{\cos(bt \sin \phi)\}(a) \\ &= \frac{1}{\pi} \frac{a}{b^2 \sin^2 \phi + a^2} \end{aligned}$$

$$\begin{aligned} \int_0^\infty e^{-at} J_0(bt) dt &= \int_0^\pi \left( \Phi(b, \phi) \sin \phi + \eta(b, \phi) \sin^2 \phi \right) d\phi \\ &= \frac{1}{\pi b} \int_0^\pi \arctan\left(\frac{b \sin \phi}{a}\right) \sin \phi d\phi + \frac{1}{\pi} \int_0^\pi \frac{a \sin^2 \phi}{b^2 \sin^2 \phi + a^2} d\phi \end{aligned}$$

$$\begin{aligned} \frac{1}{\pi} \int_0^\pi \frac{a \sin^2 \phi}{b^2 \sin^2 \phi + a^2} d\phi &= \frac{a}{\pi b^2} \int_0^\pi \frac{b^2 \sin^2 \phi}{b^2 \sin^2 \phi + a^2} d\phi \\ &= \frac{a}{\pi b^2} \int_0^\pi \frac{b^2 \sin^2 \phi + a^2 - a^2}{b^2 \sin^2 \phi + a^2} d\phi \end{aligned}$$

$$\begin{aligned}
 &= \frac{a}{\pi b^2} \int_0^\pi d\phi - \frac{a}{\pi b^2} \int_0^\pi \frac{a^2}{b^2 \sin^2 \phi + a^2} d\phi \\
 &= \frac{\pi a}{\pi b^2} - \frac{a^3}{\pi b^2} \int_0^\pi \frac{1}{b^2 \sin^2 \phi + a^2} d\phi
 \end{aligned}$$

$$\begin{aligned}
 \int_0^\pi \frac{1}{b^2 \sin^2 \phi + a^2} d\phi &\stackrel{v \rightarrow \tan(\frac{\phi}{2})}{=} \int_0^\infty \frac{1}{b^2 \left(\frac{2v}{1+v^2}\right)^2 + a^2} \frac{2dv}{1+v^2} \\
 &= 2 \int_0^\infty \frac{1+v^2}{4b^2 v^2 + a^2 (1+v^2)^2} dv \\
 &= 2 \int_0^\infty \left( \frac{1}{2a(av^2 + a + 2ibv)} + \frac{1}{2a(av^2 + a - 2ibv)} \right) dv \\
 &= \frac{1}{a^2} \int_0^\infty \left( \frac{1}{v^2 + 1 + \frac{2ibv}{a}} + \frac{1}{v^2 + 1 - \frac{2ibv}{a}} \right) dv \\
 &= \frac{1}{a^2} \int_0^\infty \left( \frac{1}{\left(v + \frac{ib}{a}\right)^2 + 1 - \left(\frac{ib}{a}\right)^2} + \frac{1}{\left(v - \frac{ib}{a}\right)^2 + 1 - \left(\frac{ib}{a}\right)^2} \right) dv \\
 &= \frac{1}{a^2} \int_0^\infty \left( \frac{1}{\left(v + \frac{ib}{a}\right)^2 + 1 + \left(\frac{b}{a}\right)^2} + \frac{1}{\left(v - \frac{ib}{a}\right)^2 + 1 + \left(\frac{b}{a}\right)^2} \right) dv \\
 &= \frac{1}{a^2} \cdot \frac{1}{\sqrt{1 + \frac{b^2}{a^2}}} \left( \arctan \left( \frac{v + \frac{ib}{a}}{\sqrt{1 + \frac{b^2}{a^2}}} \right) + \arctan \left( \frac{v - \frac{ib}{a}}{\sqrt{1 + \frac{b^2}{a^2}}} \right) \right) \Big|_0^\infty \\
 &= \frac{1}{a^2} \cdot \frac{1}{\sqrt{1 + \frac{b^2}{a^2}}} \left( \frac{\pi}{2} + \frac{\pi}{2} \right) \\
 &= \frac{\pi}{a^2} \cdot \frac{1}{\sqrt{1 + \frac{b^2}{a^2}}} = \frac{\pi}{a\sqrt{a^2 + b^2}}
 \end{aligned}$$

$$\begin{aligned}
 \therefore \frac{1}{\pi} \int_0^\pi \frac{a \sin^2 \phi}{b^2 \sin^2 \phi + a^2} d\phi &= \frac{\pi a}{\pi b^2} - \frac{a^3 \pi}{a\pi b^2 \sqrt{a^2 + b^2}} \\
 &= \frac{a}{b^2} - \frac{a^2}{b^2 \sqrt{a^2 + b^2}}
 \end{aligned}$$

$$\begin{aligned}
 \frac{1}{\pi b} \int_0^\pi \arctan \left( \frac{b \sin \phi}{a} \right) \sin \phi d\phi &= \frac{1}{\pi b} \int_0^\pi \arctan \left( \frac{b \sin \phi}{a} \right) d(-\cos \phi) \\
 &= -\frac{\cos \phi}{\pi b} \arctan \left( \frac{b \sin \phi}{a} \right) \Big|_0^\pi + \frac{1}{\pi b} \int_0^\pi \frac{\frac{b \cos \phi}{a}}{1 + \left(\frac{b \sin \phi}{a}\right)^2} \cos \phi d\phi \\
 &= \frac{1}{\pi a} \int_0^\pi \frac{\cos^2 \phi}{1 + \left(\frac{b \sin \phi}{a}\right)^2} d\phi \\
 &= \frac{a}{\pi} \int_0^\pi \frac{\cos^2 \phi}{a^2 + b^2 \sin^2 \phi} d\phi
 \end{aligned}$$

But

$$\begin{aligned} \int_0^\pi \frac{1}{b^2 \sin^2 \phi + a^2} d\phi &= \int_0^\infty \frac{1}{b^2 \left(\frac{2v}{1+v^2}\right)^2 + a^2} \frac{2dv}{1+v^2} = \frac{\pi}{a\sqrt{a^2+b^2}} \\ \implies \frac{a}{\pi} \int_0^\pi \frac{\cos^2 \phi}{a^2 + b^2 \sin^2 \phi} d\phi &= \frac{a}{\pi} \int_0^\pi \frac{1 - \sin^2 \phi}{a^2 + b^2 \sin^2 \phi} d\phi \\ &= \frac{a}{\pi} \int_0^\pi \frac{1}{b^2 \sin^2 \phi + a^2} d\phi - \frac{1}{\pi} \int_0^\pi \frac{a \sin^2 \phi}{a^2 + b^2 \sin^2 \phi} d\phi \\ &= \frac{a}{\pi} \cdot \frac{\pi}{a\sqrt{a^2+b^2}} - \left( \frac{a}{b^2} - \frac{a^2}{b^2\sqrt{a^2+b^2}} \right) \\ &= \frac{1}{\sqrt{a^2+b^2}} - \frac{a}{b^2} + \frac{a^2}{b^2\sqrt{a^2+b^2}} \end{aligned}$$

$$\begin{aligned} \implies \int_0^\infty e^{-at} J_0(bt) dt &= \int_0^\pi (\Phi(b, \phi) \sin \phi + \eta(b, \phi) \sin^2 \phi) d\phi \\ &= \frac{1}{\pi b} \int_0^\pi \arctan\left(\frac{b \sin \phi}{a}\right) \sin \phi d\phi + \frac{1}{\pi} \int_0^\pi \frac{a \sin^2 \phi}{b^2 \sin^2 \phi + a^2} d\phi \\ &= \frac{1}{\sqrt{a^2+b^2}} - \frac{a}{b^2} + \frac{a^2}{b^2\sqrt{a^2+b^2}} + \frac{a}{b^2} - \frac{a^2}{b^2\sqrt{a^2+b^2}} \\ &= \frac{1}{\sqrt{a^2+b^2}} \end{aligned}$$

Hence,

$$\int_0^\infty e^{-at} J_0(bt) dt = \frac{1}{\sqrt{a^2+b^2}}$$

**4 Evaluating**  $\int_0^\infty e^{-at} J_\alpha(bt) t^{-\alpha} dt$ ,  $\int_0^\infty J_{\frac{1}{2}}(bt) t^{-\frac{1}{2}} dt$ ,  
 $\int_0^\infty e^{-at} J_{\frac{1}{2}}(bt) t^{-\frac{1}{2}} dt$ ,  $\int_0^\infty e^{-at} J_1(bt) t^{-1} dt$ ,  $\int_0^\infty J_1(bt) t^{-1} dt$   
**and**  $\int_0^\infty J_\alpha(bt) t^{-\alpha} dt$  **with the new integral representation**

**4.1**  $\int_0^\infty e^{-at} J_\alpha(bt) t^{-\alpha} dt$

It has been established previously that

$$J_\alpha(x) = \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \left(2 \sin^{2\alpha-1} \vartheta \sin \phi [\sin(x \cos \vartheta \sin \phi) + x \cos \vartheta \sin \phi \cos(x \cos \vartheta \sin \phi)]\right) d\phi d\vartheta$$

$$\implies J_\alpha(bt) = \frac{1}{\pi b t \Gamma(\alpha)} \left(\frac{bt}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \left[2 \sin^{2\alpha-1} \vartheta \sin \phi \left(\sin(bt \cos \vartheta \sin \phi)\right)\right]$$

$$\begin{aligned}
 & + bt \cos \vartheta \sin \phi \cos (bt \cos \vartheta \sin \phi) \Big] \Big) d\phi d\vartheta \\
 \implies \int_0^\infty e^{-at} J_\alpha(bt) t^{-\alpha} dt &= \int_0^\infty \frac{1}{\pi b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha t^{-1} e^{-at} \left( \int_0^{\frac{\pi}{2}} \int_0^\pi \left[ 2 \sin^{2\alpha-1} \vartheta \sin \phi [\sin (bt \cos \vartheta \sin \phi) \right. \right. \\
 & \left. \left. + bt \cos \vartheta \sin \phi \cos (bt \cos \vartheta \sin \phi) \right] d\phi d\vartheta dt \right) \\
 &= \frac{1}{\pi b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \int_0^\infty t^{-1} e^{-at} \left( \left[ 2 \sin^{2\alpha-1} \vartheta \sin \phi [\sin (bt \cos \vartheta \sin \phi) \right. \right. \\
 & \left. \left. + bt \cos \vartheta \sin \phi \cos (bt \cos \vartheta \sin \phi) \right] \right) dt d\phi d\vartheta \Big)
 \end{aligned}$$

From our previous calculation on Lipschitz's result,

$$\int_0^\pi \left( \int_0^\infty e^{-at} \frac{\sin (bt \sin \phi) \sin \phi}{\pi bt} dt + \int_0^\infty \frac{e^{-at}}{\pi} \sin^2 \phi \cos (bt \sin \phi) dt \right) d\phi = \frac{1}{\sqrt{a^2 + b^2}} \quad (4)$$

Replacing  $b$  with  $b \cos \vartheta$  in (4) above,

$$\begin{aligned}
 \implies \int_0^\pi \int_0^\infty \frac{1}{b \pi \cos \vartheta} t^{-1} e^{-at} \left( \sin \phi [\sin (bt \cos \vartheta \sin \phi) + bt \cos \vartheta \sin \phi \cos (bt \cos \vartheta \sin \phi)] \right) dt d\phi &= \frac{1}{\sqrt{b^2 \cos^2 \vartheta + a^2}} \\
 \implies \int_0^\pi \int_0^\infty \frac{1}{b \pi} t^{-1} e^{-at} \left( \sin \phi [\sin (bt \cos \vartheta \sin \phi) + bt \cos \vartheta \sin \phi \cos (bt \cos \vartheta \sin \phi)] \right) dt d\phi &= \frac{\cos \vartheta}{\sqrt{b^2 \cos^2 \vartheta + a^2}}
 \end{aligned}$$

$$\begin{aligned}
 \int_0^\infty e^{-at} J_\alpha(bt) t^{-\alpha} dt &= \frac{1}{\Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \frac{2 \sin^{2\alpha-1} \vartheta \cos \vartheta}{\sqrt{b^2 \cos^2 \vartheta + a^2}} d\vartheta \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \frac{2 \sin^{2\alpha-1} \vartheta}{\sqrt{b^2 + a^2 - b^2 \sin^2 \vartheta}} d(\sin \vartheta) \\
 &= \frac{1}{\Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \int_0^1 \frac{2u^{2\alpha-1}}{\sqrt{b^2 + a^2 - b^2 u^2}} du \\
 &= \frac{1}{\Gamma(\alpha) \sqrt{b^2 + a^2}} \left(\frac{b}{2}\right)^\alpha \int_0^1 \frac{2u^{2\alpha-1}}{\sqrt{1 - \frac{b^2 u^2}{b^2 + a^2}}} du \tag{5} \\
 &= \frac{1}{b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \left(\frac{\sqrt{b^2 + a^2}}{b}\right)^{2\alpha-1} \int_0^{\frac{b}{\sqrt{b^2 + a^2}}} \frac{2u^{2\alpha-1}}{\sqrt{1 - u^2}} du \\
 &= \frac{1}{2b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \left(\frac{\sqrt{b^2 + a^2}}{b}\right)^{2\alpha-1} \int_0^{\frac{b}{\sqrt{b^2 + a^2}}} \frac{2u^{\alpha-1}}{\sqrt{1 - u}} du \\
 &= \frac{1}{b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \left(\frac{\sqrt{b^2 + a^2}}{b}\right)^{2\alpha-1} \int_0^{\frac{b}{\sqrt{b^2 + a^2}}} u^{\alpha-1} (1 - u)^{-\frac{1}{2}} du \\
 &= \frac{1}{b \Gamma(\alpha)} \left(\frac{b}{2}\right)^\alpha \left(\frac{\sqrt{b^2 + a^2}}{b}\right)^{2\alpha-1} B\left(\frac{b}{\sqrt{b^2 + a^2}}; \alpha, \frac{1}{2}\right) \\
 &= \frac{(b^2 + a^2)^\alpha}{2^\alpha b^{\alpha+1} \Gamma(\alpha)} \cdot \frac{b}{\sqrt{b^2 + a^2}} B\left(\frac{b}{\sqrt{b^2 + a^2}}; \alpha, \frac{1}{2}\right) \quad \text{valid for all } a \geq 0, \alpha > 0, \tag{6}
 \end{aligned}$$

where  $B(x; a, b)$  is the incomplete beta function.



**4.2**  $\int_0^\infty J_{\frac{1}{2}}(bt)t^{-\frac{1}{2}} dt$

From equation (5),

$$\int_0^\infty e^{-at} J_\alpha(bt)t^{-\alpha} dt = \frac{1}{\Gamma(\alpha)\sqrt{b^2+a^2}} \left(\frac{b}{2}\right)^\alpha \int_0^1 \frac{2u^{2\alpha-1}}{\sqrt{1-\frac{b^2u^2}{b^2+a^2}}} du$$

Let  $\alpha = \frac{1}{2}$  and  $a = 0$

$$\begin{aligned} \implies \int_0^\infty J_{\frac{1}{2}}(bt)t^{-\frac{1}{2}} dt &= \frac{1}{b\Gamma(\frac{1}{2})} \left(\frac{b}{2}\right)^{\frac{1}{2}} \int_0^1 \frac{2}{\sqrt{1-u^2}} du \\ &= \frac{2}{b\sqrt{\pi}} \sqrt{\frac{b}{2}} \arcsin u \Big|_0^1 \\ &= \frac{2}{b\sqrt{\pi}} \sqrt{\frac{b}{2}} \cdot \frac{\pi}{2} \\ &= \sqrt{\frac{\pi}{2b}} \end{aligned}$$

$$\therefore \int_0^\infty J_{\frac{1}{2}}(bt)t^{-\frac{1}{2}} dt = \sqrt{\frac{\pi}{2b}} \quad \forall \Re(b) > 0$$

**4.3**  $\int_0^\infty e^{-at} J_{\frac{1}{2}}(bt)t^{-\frac{1}{2}} dt$

$$\begin{aligned} \int_0^\infty e^{-at} J_{\frac{1}{2}}(bt)t^{-\frac{1}{2}} dt &= \frac{1}{\Gamma(\frac{1}{2})\sqrt{b^2+a^2}} \left(\frac{b}{2}\right)^{\frac{1}{2}} \int_0^1 \frac{2}{\sqrt{1-\frac{b^2u^2}{b^2+a^2}}} du \\ &= \frac{2}{\sqrt{\pi}\sqrt{b^2+a^2}} \left(\frac{b}{2}\right)^{\frac{1}{2}} \frac{\sqrt{b^2+a^2}}{b} \arcsin\left(\frac{b}{\sqrt{b^2+a^2}}u\right) \Big|_0^1 \\ &= \sqrt{\frac{2}{\pi b}} \arcsin\left(\frac{b}{\sqrt{b^2+a^2}}\right) \quad \forall \Re(b) > 0, \Re(a) \geq 0 \end{aligned}$$

**4.4**  $\int_0^\infty e^{-at} J_1(bt)t^{-1} dt$

$$\begin{aligned} \int_0^\infty e^{-at} J_1(bt)t^{-1} dt &= \frac{1}{\Gamma(1)\sqrt{b^2+a^2}} \left(\frac{b}{2}\right) \int_0^1 \frac{2u}{\sqrt{1-\frac{b^2u^2}{b^2+a^2}}} du \\ &= \frac{b}{2} \int_0^1 \frac{2u}{\sqrt{b^2+a^2-b^2u^2}} du \\ &= \frac{b}{2} \int_0^1 \frac{1}{\sqrt{b^2+a^2-b^2u}} du \\ &= \frac{b}{2} \cdot \frac{2}{-b^2} \sqrt{b^2+a^2-b^2u} \Big|_0^1 \\ &= -\frac{1}{b} \left(a - \sqrt{b^2+a^2}\right) = \frac{1}{b} \left(\sqrt{b^2+a^2} - a\right) \quad \forall \Re(b) > 0, \Re(a) \geq 0 \end{aligned}$$

**4.4.1 Special result**

$$\int_0^{\infty} e^{-3t} J_1(4t) t^{-1} dt = \frac{1}{4} (\sqrt{4^2 + 3^2} - 3) = \frac{1}{2}$$

$$\implies 2 \int_0^{\infty} e^{-3t} J_1(4t) t^{-1} dt = 1$$

$$\mathbf{4.5} \quad \int_0^{\infty} J_1(bt) t^{-1} dt$$

It has been established previously that

$$\int_0^{\infty} e^{-at} J_1(bt) t^{-1} dt = \frac{1}{b} (\sqrt{b^2 + a^2} - a)$$

At  $a = 0$ ,

$$\int_0^{\infty} J_1(bt) t^{-1} dt = \frac{1}{b} (\sqrt{b^2 + 0^2} - 0) = 1$$

$$\therefore \int_0^{\infty} J_1(bt) t^{-1} dt = 1 \quad \forall \Re(b) > 0$$

$$\mathbf{4.6} \quad \int_0^{\infty} e^{-at} J_{\alpha}(bt) t^{-\alpha} dt$$

From equation (6)

$$\int_0^{\infty} e^{-at} J_{\alpha}(bt) t^{-\alpha} dt = \frac{(b^2 + a^2)^{\alpha}}{2^{\alpha} b^{\alpha+1} \Gamma(\alpha)} \cdot \frac{b}{\sqrt{b^2 + a^2}} B\left(\frac{b}{\sqrt{b^2 + a^2}}; \alpha, \frac{1}{2}\right) \quad \forall \alpha > 0, a \geq 0$$

At  $a = 0$ ,

$$\begin{aligned} \int_0^{\infty} J_{\alpha}(bt) t^{-\alpha} dt &= \frac{b^{\alpha-1}}{2^{\alpha} \Gamma(\alpha)} B\left(1; \alpha, \frac{1}{2}\right) \\ &= \frac{b^{\alpha-1}}{2^{\alpha} \Gamma(\alpha)} \beta\left(\alpha, \frac{1}{2}\right) \\ &= \frac{b^{\alpha-1}}{2^{\alpha} \Gamma(\alpha)} \cdot \frac{\Gamma(\alpha) \Gamma(\frac{1}{2})}{\Gamma(\alpha + \frac{1}{2})} \\ &= \frac{b^{\alpha-1} \sqrt{\pi}}{2^{\alpha} \Gamma(\alpha + \frac{1}{2})} \\ &= \frac{b^{\alpha-1} \sqrt{\pi}}{2^{\alpha} \left(\frac{\Gamma(2\alpha) \sqrt{\pi}}{\Gamma(\alpha) 2^{2\alpha-1}}\right)} = \frac{b^{\alpha-1}}{\Gamma(\alpha) 2^{\alpha-1}} \\ &= \frac{(2b)^{\alpha-1} \Gamma(\alpha)}{\Gamma(2\alpha)} \end{aligned}$$

$$\therefore \int_0^{\infty} J_{\alpha}(bt) t^{-\alpha} dt = \begin{cases} \frac{(2b)^{\alpha-1} \Gamma(\alpha)}{\Gamma(2\alpha)} & \text{for } \alpha > 0 \\ \frac{b^{\alpha-1} \sqrt{\pi}}{2^{\alpha} \Gamma(\alpha + \frac{1}{2})} = \frac{1}{b} & \text{for } \alpha = 0 \end{cases}$$

#### 4.6.1 Special results

$$\int_0^\infty J_2(3t)t^{-2} dt = \frac{(2 \times 3)^1 \Gamma(2)}{\Gamma(4)} = 1$$

$$\int_0^\infty J_0(bt) dt = \frac{1}{b} \quad \forall \Re(b) > 0$$

$$\int_0^\infty J_0(t) dt = 1 \quad \text{for } b = 1.$$

## 5 Integral representations for $I_\alpha(x)$ and $I_0(x)$

The Bessel functions are valid even for complex arguments  $x$ , and an important special case is that of a purely imaginary argument. In this case, the solutions to the Bessel equation are called the modified Bessel functions (or occasionally the hyperbolic Bessel functions) of the first and second kind and are defined as [5]

$$I_\alpha(x) = i^{-\alpha} J_\alpha(ix) = \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m + \alpha + 1)} \left(\frac{x}{2}\right)^{2m + \alpha} \quad \forall x \in \mathbb{C}$$

$$K_\alpha(x) = \frac{\pi I_{-\alpha}(x) - I_\alpha(x)}{2 \sin(\alpha\pi)} \quad \text{when } \alpha \text{ is not an integer.} \quad (7)$$

Now, at  $\alpha = 0$ ,

$$I_0(x) = J_0(ix) \quad (8)$$

By equation (3),

$$J_0(x) = \frac{1}{\pi x} \int_0^\pi (\sin(x \sin \phi) + x \sin \phi \cos(x \sin \phi)) \sin \phi d\phi$$

$$\begin{aligned} I_0(x) = J_0(ix) &= \frac{1}{\pi ix} \int_0^\pi (\sin(ix \sin \phi) + ix \sin \phi \cos(ix \sin \phi)) \sin \phi d\phi \\ &= \frac{1}{\pi x} \int_0^\pi (-i \sin(ix \sin \phi) + x \sin \phi \cosh(x \sin \phi)) \sin \phi d\phi \\ &= \frac{1}{\pi x} \int_0^\pi (\sinh(x \sin \phi) + x \sin \phi \cosh(x \sin \phi)) \sin \phi d\phi \end{aligned}$$

$$\implies I_0(x) = \frac{1}{\pi x} \int_0^\pi (\sinh(x \sin \phi) + x \sin \phi \cosh(x \sin \phi)) \sin \phi d\phi \quad (9)$$

By equation (1),

$$\begin{aligned} J_\alpha(ix) &= \frac{1}{\Gamma(\alpha)} \left(\frac{ix}{2}\right)^\alpha \int_0^1 J_0(ix\sqrt{1-t}) t^{\alpha-1} dt \\ \implies i^{-\alpha} J_\alpha(ix) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 J_0(ix\sqrt{1-t}) t^{\alpha-1} dt \end{aligned}$$

By equation (8),  $I_0(x) = J_0(ix)$ .

$$\begin{aligned} \implies i^{-\alpha} J_\alpha(ix) = I_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 I_0(x\sqrt{1-t}) t^{\alpha-1} dt \\ \implies I_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 I_0(x\sqrt{1-t}) t^{\alpha-1} dt \end{aligned} \quad (10)$$

Inserting (9) in (10),

$$I_\alpha(x) = \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi \left[ 2 \sin^{2\alpha-1} \vartheta \sin \phi \left[ \sinh(x \cos \vartheta \sin \phi) + x \cos \vartheta \sin \phi \cosh(x \cos \vartheta \sin \phi) \right] \right] d\phi d\vartheta$$

## 6 Simplifications of the new integral representations

$$\begin{aligned} J_0(x) &= \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} \\ &= \frac{1}{\pi x} \int_0^\pi (\sin(x \sin \phi) + x \sin \phi \cos(x \sin \phi)) \sin \phi d\phi \\ &= \frac{1}{\pi x} \int_0^\pi \left( \sin(x \sin \phi) + x \frac{\partial}{\partial x} (\sin(x \sin \phi)) \right) \sin \phi d\phi \\ &= \frac{1}{\pi x} \int_0^\pi \sin \phi \frac{\partial}{\partial x} (x \sin(x \sin \phi)) d\phi \\ \implies J_0(x) &= \frac{1}{\pi x} \int_0^\pi \sin \phi \frac{\partial}{\partial x} (x \sin(x \sin \phi)) d\phi \end{aligned}$$

Since

$$\begin{aligned} I_0(x) &= \frac{1}{\pi x} \int_0^\pi (\sinh(x \sin \phi) + x \sin \phi \cosh(x \sin \phi)) \sin \phi d\phi \\ \implies I_0(x) &= \frac{1}{\pi x} \int_0^\pi \sin \phi \frac{\partial}{\partial x} (x \sinh(x \sin \phi)) d\phi \end{aligned}$$

$$I_\alpha(x) = \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} 2 \sin^{2\alpha-1} \vartheta \cos \vartheta I_0(x \cos \vartheta) d\vartheta$$

$$\begin{aligned} \implies I_\alpha(x) &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi 2 \sin^{2\alpha-1} \vartheta \cos \vartheta \frac{\sin \phi}{\cos \vartheta} \frac{\partial}{\partial (x \cos \vartheta)} (x \cos \vartheta \sinh(x \cos \vartheta \sin \phi)) d\phi d\vartheta \\ &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi 2 \sin^{2\alpha-1} \vartheta \sin \phi \frac{\partial}{\partial (x \cos \vartheta)} (x \cos \vartheta \sinh(x \cos \vartheta \sin \phi)) d\phi d\vartheta \end{aligned}$$

Similarly,

$$J_\alpha(x) = \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi 2 \sin^{2\alpha-1} \vartheta \sin \phi \frac{\partial}{\partial (x \cos \vartheta)} (x \cos \vartheta \sin(x \cos \vartheta \sin \phi)) d\phi d\vartheta$$

## 7 An expression for $K_{\frac{1}{2}}(x) \forall \Re(x) > 0$ with the new integral representations

$$I_\alpha(x) = \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} 2 \sin^{2\alpha-1} \vartheta \cos \vartheta I_0(x \cos \vartheta) d\vartheta$$

$$\begin{aligned} I_{\frac{1}{2}}(x) &= \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} 2 \sin^{2(\frac{1}{2})-1} \vartheta \cos \vartheta I_0(x \cos \vartheta) d\vartheta \\ &= \frac{2}{\sqrt{\pi}} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} \cos \vartheta I_0(x \cos \vartheta) d\vartheta \end{aligned}$$

$$J_{-\frac{1}{2}}(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m - \frac{1}{2} + 1)} \left(\frac{x}{2}\right)^{2m - \frac{1}{2}}$$

$$\begin{aligned}
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \sum_{m=1}^{\infty} \frac{(-1)^m}{(m+1)! \Gamma(m-\frac{1}{2}+1)} \left(\frac{x}{2}\right)^{2m-\frac{1}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \sum_{m=0}^{\infty} \frac{(-1)^{m+1}}{(m+1)! \Gamma(m+\frac{1}{2}+1)} \left(\frac{x}{2}\right)^{2m+\frac{3}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \sum_{m=0}^{\infty} \frac{(-1)^m \Gamma(m+1) \Gamma(\frac{1}{2})}{\Gamma(m+2) \Gamma(m+1) \Gamma(m+\frac{1}{2}+1)} \left(\frac{x}{2}\right)^{2m+\frac{3}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \sum_{m=0}^{\infty} \frac{(-1)^m \Gamma(m+1) \Gamma(\frac{1}{2})}{m! \Gamma(m+2) \Gamma(m+\frac{1}{2}+1)} \left(\frac{x}{2}\right)^{2m+\frac{3}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+2)} \beta\left(m+1, \frac{1}{2}\right) \left(\frac{x}{2}\right)^{2m+\frac{3}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+2)} \int_0^1 t^m (1-t)^{\frac{1}{2}-1} dt \left(\frac{x}{2}\right)^{2m+\frac{3}{2}} \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+2)} t^m \left(\frac{x}{2}\right)^{2m+1} (1-t)^{-\frac{1}{2}} dt \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+2)} \left(\frac{x\sqrt{t}}{2}\right)^{2m+1} t^{-\frac{1}{2}} (1-t)^{-\frac{1}{2}} dt \\
 \\
 \implies J_{-\frac{1}{2}}(x) &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 J_1(x\sqrt{t}) t^{-\frac{1}{2}} (1-t)^{-\frac{1}{2}} dt \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 J_1(x\sqrt{1-t}) t^{\frac{1}{2}-1} (1-t)^{-\frac{1}{2}} dt
 \end{aligned}$$

Since,  $I_{\alpha}(x) = i^{-\alpha} J_{\alpha}(ix)$ ,

$$\implies I_{-\frac{1}{2}}(x) = i^{\frac{1}{2}} J_{\frac{1}{2}}(ix)$$

$$I_1(x) = i^{-1} J_1(ix) \implies J_1(ix) = iI_1(x) \implies iJ_1(ix) = -I_1(x)$$

$$\begin{aligned}
 \implies i^{\frac{1}{2}} J_{\frac{1}{2}}(ix) &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} - \frac{i}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 J_1(ix\sqrt{1-t}) t^{\frac{1}{2}-1} (1-t)^{-\frac{1}{2}} dt \\
 \implies I_{-\frac{1}{2}}(x) &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^1 I_1(x\sqrt{1-t}) t^{-\frac{1}{2}} (1-t)^{-\frac{1}{2}} dt
 \end{aligned}$$

$$\begin{aligned}
 I_{-\frac{1}{2}}(x) &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \frac{1}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} 2 \sin^{2(\frac{1}{2})-1} \vartheta \cos \vartheta \frac{1}{\cos \vartheta} I_1(x \cos \vartheta) d\vartheta \\
 &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \frac{2}{\Gamma(\frac{1}{2})} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} I_1(x \cos \vartheta) d\vartheta
 \end{aligned}$$

$$\begin{aligned}
 \int_0^{\frac{\pi}{2}} \cos \vartheta I_0(x \cos \vartheta) d\vartheta &= \int_0^1 I_0(x\sqrt{1-u^2}) du \\
 &= \frac{1}{2} \int_0^1 u^{-\frac{1}{2}} I_0(x\sqrt{1-u}) du
 \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \int_0^1 (1-u)^{-\frac{1}{2}} I_0(x\sqrt{u}) \, du \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+1)} \int_0^1 (1-u)^{-\frac{1}{2}} \left(\frac{x\sqrt{u}}{2}\right)^{2m} \, du \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} \int_0^1 u^m (1-u)^{-\frac{1}{2}} \, du \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} \beta\left(m+1, \frac{1}{2}\right) \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} \frac{\Gamma(m+1) \Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})} \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m!} \left(\frac{x}{2}\right)^{2m} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})} \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{\Gamma(m+1)} \left(\frac{x}{2}\right)^{2m} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})}
\end{aligned}$$

By Legendre's duplication formula,

$$\begin{aligned}
\frac{2^{2m+1}}{\Gamma(2m+2)} &= \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2}) \Gamma(m+1)} \\
\Rightarrow \int_0^{\frac{\pi}{2}} \cos \vartheta I_0(x \cos \vartheta) \, d\vartheta &= \frac{1}{2} \sum_{m=0}^{\infty} \frac{2^{2m+1}}{\Gamma(2m+2)} \left(\frac{x}{2}\right)^{2m} \\
&= \sum_{m=0}^{\infty} \frac{x^{2m}}{\Gamma(2m+2)} = \sum_{m=0}^{\infty} \frac{x^{2m}}{(2m+1)!} \\
&= \frac{1}{x} \sum_{m=0}^{\infty} \frac{x^{2m+1}}{(2m+1)!} \\
&= \frac{\sinh(x)}{x}
\end{aligned}$$

$$\begin{aligned}
\Rightarrow I_{\frac{1}{2}}(x) &= \frac{2}{\sqrt{\pi}} \left(\frac{x}{2}\right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} \cos \vartheta I_0(x \cos \vartheta) \, d\vartheta \\
&= \sqrt{\frac{2x}{\pi}} \frac{\sinh(x)}{x} \\
&= \sqrt{\frac{2}{\pi x}} \sinh(x)
\end{aligned}$$

$$\begin{aligned}
\int_0^{\frac{\pi}{2}} I_1(x \cos \vartheta) \, d\vartheta &= \int_0^1 \frac{I_1(x\sqrt{1-u^2})}{\sqrt{1-u^2}} \, du \\
&= \frac{1}{2} \int_0^1 u^{-\frac{1}{2}} (1-u)^{-\frac{1}{2}} I_1(x\sqrt{1-u}) \, du \\
&= \frac{1}{2} \int_0^1 u^{-\frac{1}{2}} (1-u)^{-\frac{1}{2}} I_1(x\sqrt{u}) \, du \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+2)} \int_0^1 u^{-\frac{1}{2}} (1-u)^{-\frac{1}{2}} \left(\frac{x\sqrt{u}}{2}\right)^{2m+1} \, du
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+2)} \left(\frac{x}{2}\right)^{2m+1} \int_0^1 u^m (1-u)^{-\frac{1}{2}} du \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m+2)} \left(\frac{x}{2}\right)^{2m+1} \beta\left(m+1, \frac{1}{2}\right) \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m!(m+1)\Gamma(m+1)} \left(\frac{x}{2}\right)^{2m+1} \frac{\Gamma(m+1)\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})} \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{m!(m+1)} \left(\frac{x}{2}\right)^{2m+1} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})} \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{1}{\Gamma(m+2)} \left(\frac{x}{2}\right)^{2m+1} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{3}{2})} \\
&= \frac{1}{2} \sum_{m=0}^{\infty} \frac{m+\frac{3}{2}}{\Gamma(m+2)} \left(\frac{x}{2}\right)^{2m+1} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{5}{2})} \\
&= \frac{1}{4} \sum_{m=0}^{\infty} \frac{2m+3}{\Gamma(m+2)} \left(\frac{x}{2}\right)^{2m+1} \frac{\Gamma(\frac{1}{2})}{\Gamma(m+\frac{5}{2})}
\end{aligned}$$

$$\begin{aligned}
\Rightarrow \int_0^{\frac{\pi}{2}} \cos \vartheta I_0(x \cos \vartheta) d\vartheta &= \frac{1}{4} \sum_{m=0}^{\infty} \frac{2^{2m+3}}{\Gamma(2m+4)} (2m+3) \left(\frac{x}{2}\right)^{2m+1} \\
&= 2 \cdot \frac{x}{2} \sum_{m=0}^{\infty} \frac{x^{2m} (2m+3)}{(2m+3)\Gamma(2m+3)} = x \sum_{m=0}^{\infty} \frac{x^{2m}}{(2m+2)!} \\
&= \frac{1}{x} \sum_{m=0}^{\infty} \frac{x^{2m+2}}{(2m+2)!} \\
&= \frac{\cosh(x) - 1}{x}
\end{aligned}$$

$$\begin{aligned}
I_{-\frac{1}{2}}(x) &= \frac{1}{\Gamma(\frac{1}{2})} \sqrt{\frac{2}{x}} + \sqrt{\frac{2x}{\pi}} \left(\frac{\cosh(x) - 1}{x}\right) \\
&= \sqrt{\frac{2}{\pi x}} + \sqrt{\frac{2x}{\pi}} \left(\frac{\cosh(x) - 1}{x}\right) \\
&= \sqrt{\frac{2}{\pi x}} \cosh(x)
\end{aligned}$$

$$\begin{aligned}
I_{-\frac{1}{2}}(x) - I_{\frac{1}{2}}(x) &= \sqrt{\frac{2}{\pi x}} \cosh(x) - \sqrt{\frac{2}{\pi x}} \sinh(x) \\
&= \sqrt{\frac{2}{\pi x}} e^{-x}
\end{aligned}$$

At  $\alpha = \frac{1}{2}$  in equation (7),

$$\begin{aligned}
K_{\frac{1}{2}}(x) &= \frac{\pi}{2} \frac{I_{-\frac{1}{2}}(x) - I_{\frac{1}{2}}(x)}{\sin(\frac{\pi}{2})} \\
&= \frac{\pi}{2} \left(I_{-\frac{1}{2}}(x) - I_{\frac{1}{2}}(x)\right)
\end{aligned}$$

$$\begin{aligned}
&= \frac{\pi}{2} \cdot \sqrt{\frac{2}{\pi x}} e^{-x} = \sqrt{\frac{\pi}{2x}} e^{-x} \\
\therefore K_{\frac{1}{2}}(x) &= \sqrt{\frac{\pi}{2x}} e^{-x} \tag{11}
\end{aligned}$$

## 8 Deriving another integral representation for $K_{\alpha}(xz)$ directly from Basset's formula

By Basset's formula [3],

$$K_{\alpha}(xz) = \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2z)^{\alpha}}{x^{\alpha} \Gamma(\frac{1}{2})} \int_0^{\infty} \frac{\cos(xu)}{(u^2 + z^2)^{\alpha + \frac{1}{2}}} du$$

By elementary substitution, Basset's formula can be rewritten to  $K_{\alpha}(z)$ .

$$\begin{aligned}
K_{\alpha}(xz) &= \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2z)^{\alpha}}{x^{\alpha} \Gamma(\frac{1}{2})} \int_0^{\infty} \frac{\cos(xu)}{(u^2 + z^2)^{\alpha + \frac{1}{2}}} du \\
&= \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2z)^{\alpha}}{x^{\alpha} \Gamma(\frac{1}{2})} z \int_0^{\infty} \frac{\cos(xzu)}{(z^2 u^2 + z^2)^{\alpha + \frac{1}{2}}} du \\
&= \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2z)^{\alpha}}{(z^{2\alpha} \cdot z) x^{\alpha} \Gamma(\frac{1}{2})} z \int_0^{\infty} \frac{\cos(xzu)}{(u^2 + 1)^{\alpha + \frac{1}{2}}} du \\
&= \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2)^{\alpha}}{z^{\alpha} x^{\alpha} \Gamma(\frac{1}{2})} \int_0^{\infty} \frac{\cos(xzu)}{(u^2 + 1)^{\alpha + \frac{1}{2}}} du \\
\Rightarrow K_{\alpha}(z) &= \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2)^{\alpha}}{z^{\alpha} \Gamma(\frac{1}{2})} \int_0^{\infty} \frac{\cos(zu)}{(u^2 + 1)^{\alpha + \frac{1}{2}}} du
\end{aligned}$$

Consider the following integral

$$H(\beta, p, R) = \frac{2\Gamma(p)}{\sqrt{\pi}} \int_0^{\infty} \frac{\cos(2Rx)}{(\beta^2 + x^2)^p} dx$$

$$\begin{aligned}
H(\beta, p, R) &= \frac{2\Gamma(p)}{\sqrt{\pi}} \int_0^{\infty} \frac{\cos(2Rx)}{(\beta^2 + x^2)^p} dx \\
&= \frac{2}{\sqrt{\pi}} \int_0^{\infty} \int_0^{\infty} t^{p-1} e^{-(\beta^2 + x^2)t} \cos(2Rx) dt dx \\
&= \frac{2}{\sqrt{\pi}} \int_0^{\infty} \int_0^{\infty} t^{p-1} e^{-(\beta^2 + x^2)t} \cos(2Rx) dx dt \\
&= \frac{2}{\sqrt{\pi}} \int_0^{\infty} t^{p-1} e^{-\beta^2 t} \int_0^{\infty} e^{-x^2 t} \cos(2Rx) dx dt
\end{aligned}$$

Let

$$f(R, t) = \int_0^{\infty} e^{-x^2 t} \cos(2Rx) dx$$

$$\begin{aligned}
\frac{\partial f(R, t)}{\partial R} &= -2 \int_0^{\infty} x e^{-x^2 t} \sin(2Rx) dx \\
&= \frac{1}{t} \int_0^{\infty} \sin(2Rx) d(e^{-x^2 t})
\end{aligned}$$



$$= -\frac{2R}{t} \int_0^\infty e^{-x^2t} \cos(2Rx) dx = -\frac{2R}{t} f(R,t)$$

$$\frac{\partial f(R,t)}{\partial R} = -\frac{2R}{t} f(R,t)$$

$$\implies \frac{1}{f(R,t)} \frac{\partial f(R,t)}{\partial R} = -\frac{2R}{t}$$

$$\implies \frac{\partial (\ln(f(R,t)))}{\partial R} = -\frac{2R}{t}$$

$$\implies \partial (\ln(f(R,t))) = -\frac{2R}{t} \partial R$$

$$\implies \int d(\ln(f(R,t))) = -\int \frac{2R}{t} dR$$

$$\implies \ln(f(R,t)) = -\frac{R^2}{t} + C(t)$$

$$\implies f(R,t) = A(t)e^{-\frac{R^2}{t}}$$

$$\implies f(R,t) = \int_0^\infty e^{-x^2t} \cos(2Rx) dx = A(t)e^{-\frac{R^2}{t}}$$

$$\therefore f(0,t) = \int_0^\infty e^{-x^2t} dx = A(t).$$

$$\begin{aligned} \int_0^\infty e^{-x^2t} dx &\stackrel{x \rightarrow \sqrt{u}}{=} \frac{1}{2} \int_0^\infty u^{-\frac{1}{2}} e^{-ut} du \\ &= \frac{1}{2} \mathcal{L}\left(u^{-\frac{1}{2}}\right)(t) = \frac{\Gamma\left(\frac{1}{2}\right)}{2t^{\frac{1}{2}}} = \frac{1}{2} \sqrt{\frac{\pi}{t}} \end{aligned}$$

$$\implies f(R,t) = \int_0^\infty e^{-x^2t} \cos(2Rx) dx = \frac{1}{2} \sqrt{\frac{\pi}{t}} e^{-\frac{R^2}{t}}$$

$$\implies H(\beta, p, R) = \frac{2}{\sqrt{\pi}} \int_0^\infty t^{p-1} e^{-\beta^2t} \int_0^\infty e^{-x^2t} \cos(2Rx) dx dt$$

$$= \frac{2}{\sqrt{\pi}} \int_0^\infty t^{p-1} e^{-\beta^2t} \left(\frac{1}{2} \sqrt{\frac{\pi}{t}} e^{-\frac{R^2}{t}}\right) dt$$

$$= \int_0^\infty t^{p-\frac{3}{2}} e^{-\left(\beta^2t + \frac{R^2}{t}\right)} dt$$

$$\implies H(\beta, p, R) = \int_0^\infty t^{p-\frac{3}{2}} e^{-\left(\beta^2t + \frac{R^2}{t}\right)} dt = \frac{2\Gamma(p)}{\sqrt{\pi}} \int_0^\infty \frac{\cos(2Rt)}{(\beta^2 + t^2)^p} dt$$

Replacing  $\beta^2$  with  $a$ ,  $R^2$  with  $z$ ,

$$\int_0^\infty t^{p-\frac{3}{2}} e^{-\left(at + \frac{z}{t}\right)} dt = \frac{2\Gamma(p)}{\sqrt{\pi}} \int_0^\infty \frac{\cos(2\sqrt{zt})}{(a^2 + t^2)^p} dt$$

Replacing  $p$  with  $\frac{1}{2} + \alpha$ ,  $z$  with  $\frac{z^2}{4}$ ,  $a$  with 1,

$$\int_0^\infty t^{\alpha-1} e^{-\left(t + \frac{z}{4t}\right)} dt = \frac{2\Gamma\left(\frac{1}{2} + \alpha\right)}{\sqrt{\pi}} \int_0^\infty \frac{\cos(zt)}{(1+t^2)^{\frac{1}{2} + \alpha}} dt$$

Multiplying both sides by  $\frac{2^{\alpha-1}}{z^\alpha}$ ,

$$\frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2)^\alpha}{z^\alpha \Gamma(\frac{1}{2})} \int_0^\infty \frac{\cos(zt)}{(t^2 + 1)^{\alpha + \frac{1}{2}}} dt = \frac{1}{z^\alpha} \int_0^\infty (2t)^{\alpha-1} e^{-\left(t + \frac{t^2}{4}\right)} dt$$

Thus if

$$K_\alpha(z) = \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2)^\alpha}{z^\alpha \Gamma(\frac{1}{2})} \int_0^\infty \frac{\cos(zt)}{(t^2 + 1)^{\alpha + \frac{1}{2}}} dt$$

Then

$$K_\alpha(z) = \frac{1}{z^\alpha} \int_0^\infty (2t)^{\alpha-1} e^{-\left(t + \frac{t^2}{4}\right)} dt,$$

$$\implies K_\alpha(z) = \frac{\Gamma(\alpha + \frac{1}{2}) \cdot (2)^\alpha}{z^\alpha \Gamma(\frac{1}{2})} \int_0^\infty \frac{\cos(zt)}{(t^2 + 1)^{\alpha + \frac{1}{2}}} dt = \frac{1}{z^\alpha} \int_0^\infty (2t)^{\alpha-1} e^{-\left(t + \frac{t^2}{4}\right)} dt$$

## 9 Another variant of Hardy’s integral in relation to partial differential equations

Prof. Godfrey Harold Hardy, FRS gave the following integral representation [4]

$$\int_0^\infty \sin\left(au + \frac{b}{u}\right) \frac{du}{u} = \pi J_0\{2\sqrt{ab}\}$$

On the transformation  $u \rightarrow u^2$ , this is equivalent to

$$\int_0^\infty \sin\left(au^2 + \frac{b}{u^2}\right) \frac{du}{u} = \frac{\pi}{2} J_0\{2\sqrt{ab}\}$$

Sequel to the above, a similar integral representation for  $K_{\frac{1}{2}}(z)$  can be proposed in what follows

**Proposition.**

$$\int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du = \left(\frac{b}{2a}\right)^{\frac{1}{4}} K_{\frac{1}{2}}\{2\sqrt{ab}\}$$

**Proof.** We can take the partial derivatives of the integrand wrt.  $a$  or  $b$  by differentiating under the integral sign since the partial derivatives are continuous in the limits of the integral.

$$\frac{\partial}{\partial a} \left( \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du \right) = \int_0^\infty u^2 \cos\left(au^2 - \frac{b}{u^2}\right) du$$

$$\begin{aligned} \frac{\partial^2}{\partial a \partial b} \left( \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du \right) &= \int_0^\infty -\frac{1}{u^2} \cdot u^2 - \sin\left(au^2 - \frac{b}{u^2}\right) du \\ &= \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du \end{aligned}$$

$$\frac{\partial^2}{\partial a \partial b} \left( \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du \right) = \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du$$

Let

$$\int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du = F(a, b)$$

$$\implies \frac{\partial^2}{\partial a \partial b} F(a, b) = F(a, b)$$

Let  $F(a, b)$  be the product of two functions  $A(a), B(b)$ ,

$$\implies F(a, b) = A(a)B(b)$$

$$\implies A'(a)B'(b) = A(a)B(b)$$

$$\implies \frac{A'(a)}{A(a)} = \frac{B(b)}{B'(b)}$$

Let

$$\frac{A'(a)}{A(a)} = \frac{B(b)}{B'(b)} = \lambda(a, b)$$

$$\implies \frac{A'(a)}{A(a)} = \frac{B(b)}{B'(b)} = \lambda(a, b)$$

$$\implies \frac{A'(a)}{A(a)} = \lambda(a, b), \quad \frac{B(b)}{B'(b)} = \lambda(a, b)$$

$$\implies \frac{A'(a)}{A(a)} = \lambda(a, b), \quad \frac{B'(b)}{B(b)} = \frac{1}{\lambda(a, b)}$$

$$\implies \frac{d(A(a))}{A(a)} = \lambda(a, b) da, \quad \frac{d(B(b))}{B(b)} = \frac{1}{\lambda(a, b)} db$$

$$\implies \int \frac{d(A(a))}{A(a)} = \int \lambda(a, b) da, \quad \int \frac{d(B(b))}{B(b)} = \int \frac{1}{\lambda(a, b)} db$$

$$\implies \ln(A(a)) = \int \lambda(a, b) da + c_1, \quad \ln(B(b)) = \int \frac{db}{\lambda(a, b)} + c_2$$

$$\implies A(a) = C_1 e^{\int \lambda(a, b) da}, \quad B(b) = C_2 e^{\int \frac{db}{\lambda(a, b)}}$$

Since  $F(a, b) = A(a)B(b)$

$$\implies F(a, b) = C_1 C_2 e^{\int \lambda(a, b) da + \int \frac{db}{\lambda(a, b)}} = C_3 \exp\left(\int \lambda(a, b) da + \int \frac{db}{\lambda(a, b)}\right)$$

$$\begin{aligned} F(a, 0) &= C_3 \exp\left(\int \lambda(a, 0) da + \int \frac{db}{\lambda(a, 0)}\right) \\ &= \int_0^\infty \sin\left(au^2 - \frac{0}{u^2}\right) du = \int_0^\infty \sin(au^2) du \\ &= \frac{1}{2\sqrt{a}} \sqrt{\frac{\pi}{2}} \exp(0) \end{aligned}$$

$$\implies C_3 = \frac{1}{2\sqrt{a}} \sqrt{\frac{\pi}{2}}, \quad \int \lambda(a, 0) da + \int \frac{db}{\lambda(a, 0)} = 0$$

Let  $\lambda(a, 0) = f(a)$

$$\begin{aligned} \int f(a) da + \int \frac{db}{f(a)} &= 0 \\ \int f(a) da &= - \int \frac{db}{f(a)} = - \frac{b}{f(a)} + \text{constant} \end{aligned}$$

Taking the derivative at  $a$ ,

$$\begin{aligned}\frac{d}{da} \left( \int f(a) da \right) &= -\frac{d}{da} \left( \frac{b}{f(a)} \right) \\ \Rightarrow f(a) &= -b \cdot \frac{-f'(a)}{f(a)^2} = \frac{bf'(a)}{f(a)^2} \\ \Rightarrow f'(a) &= \frac{f(a)^3}{b} \\ \Rightarrow \frac{d(f(a))}{f(a)^3} &= \frac{da}{b} \\ \Rightarrow \int \frac{d(f(a))}{f(a)^3} &= \int \frac{da}{b}\end{aligned}$$

For  $\lambda_1, \gamma_1$  constant,

$$\begin{aligned}\Rightarrow -\frac{1}{2f(a)^2} &= \frac{a}{b} + \lambda_1 \\ \Rightarrow \frac{1}{f(a)^2} &= -\frac{2a}{b} + \lambda_1 = \frac{b\gamma_1 - 2a}{b} \\ \Rightarrow f(a) &= \sqrt{\frac{b}{b\gamma_1 - 2a}}\end{aligned}$$

$$\begin{aligned}F(0, b) &= C_3 \exp \left( \int \lambda(0, b) da + \int \frac{db}{\lambda(0, b)} \right) \\ &= \int_0^\infty \sin \left( (0)u^2 - \frac{b}{u^2} \right) du = \int_0^\infty \sin \left( -\frac{b}{u^2} \right) du \\ &= -\sqrt{b} \sqrt{\frac{\pi}{2}} \exp(0) \\ \Rightarrow C_3 &= -\sqrt{b} \sqrt{\frac{\pi}{2}}, \quad \int \lambda(0, b) da + \int \frac{db}{\lambda(0, b)} = 0\end{aligned}$$

Let  $\lambda(0, b) = f(b)$

$$\begin{aligned}\int f(b) da + \int \frac{db}{f(b)} &= 0 \\ af(b) + \int \frac{db}{f(b)} &= \text{constant} \\ af(b) &= -\int \frac{db}{f(b)} + \text{constant}\end{aligned}$$

Taking the derivative at  $b$ ,

$$\begin{aligned}\frac{d}{db} (af(b)) &= \frac{d}{db} \left( -\int \frac{db}{f(b)} \right) \\ \Rightarrow af'(b) &= -\frac{1}{f(b)} \\ \Rightarrow f'(b)f(b) &= \frac{-1}{a} \\ \Rightarrow f(b) \frac{d(f(b))}{db} &= \frac{-1}{a}\end{aligned}$$

$$\begin{aligned} \implies \frac{1}{2}d(f(b)^2) &= \frac{-db}{a} \\ \implies \frac{1}{2}\int d(f(b)^2) &= \int \frac{-db}{a} \end{aligned}$$

For  $\lambda_2, \gamma_2$  constant,

$$\begin{aligned} \implies \frac{f(b)^2}{2} &= -\frac{b}{a} + \lambda_2 \\ \implies f(b) &= \sqrt{\gamma_2 - \frac{2b}{a}} = \sqrt{\frac{a\gamma_2 - 2b}{a}} \end{aligned}$$

Since  $f(b) = \lambda(0, b)$  and  $f(a) = \lambda(a, 0)$ ,

$$\implies \lambda(0, b) = \sqrt{\frac{a\gamma_2 - 2b}{a}} = \phi_1 \sqrt{\phi_2 - 2b} \text{ and } \lambda(a, 0) = \sqrt{\frac{b}{b\gamma_1 - 2a}} = \frac{\phi_3}{\sqrt{\phi_4 - 2a}}$$

$\lambda(a, b)$  can be written from  $\lambda(0, b)$  and  $\lambda(a, 0)$

$$\implies \lambda(a, b) = \phi \sqrt{\frac{\phi_2 - 2b}{\phi_4 - 2a}} \ni \phi = \frac{\phi_3}{\sqrt{\phi_2}} \text{ or } \phi = \phi_1 \sqrt{\phi_4}.$$

$\phi \neq 0$  since we are looking for non-trivial solutions of the PDE.

$$\implies \lambda(a, b) = \phi \sqrt{\frac{\phi_2 - 2b}{\phi_4 - 2a}}, \quad \phi \neq 0$$

$$\begin{aligned} F(a, b) &= C_3 \exp\left(\phi \int \sqrt{\frac{\phi_2 - 2b}{\phi_4 - 2a}} da + \int \frac{1}{\phi} \sqrt{\frac{\phi_4 - 2b}{\phi_2 - 2a}} db\right) \\ &= C_3 \exp\left(-\phi \sqrt{(\phi_2 - 2b)(\phi_4 - 2a)} - \frac{1}{\phi} \sqrt{(\phi_2 - 2b)(\phi_4 - 2a)}\right) \\ &= C_3 \exp\left(-\left(\phi + \frac{1}{\phi}\right) \sqrt{(\phi_2 - 2b)(\phi_4 - 2a)}\right) \end{aligned}$$

Since  $\phi \neq 0, \phi + \frac{1}{\phi} \neq 0$ . So, let  $\delta = \phi + \frac{1}{\phi} \ni \delta \neq 0$ .

$$\implies F(a, b) = C_3 \exp\left(-\delta \sqrt{(\phi_2 - 2b)(\phi_4 - 2a)}\right)$$

$$\begin{aligned} F(a, 0) &= C_3 \exp\left(-\delta \sqrt{(\phi_2)(\phi_4 - 2a)}\right) \\ &= \frac{1}{2\sqrt{a}} \sqrt{\frac{\pi}{2}} \exp(0) \\ \implies C_3 &= \frac{1}{2\sqrt{a}}, \quad \phi_2 = 0. \end{aligned}$$

$\phi_4 \neq 2a$ , since  $\phi_4$  is an arbitrary constant, it cannot be a function of  $a$ .

$$F(0, b) = C_3 \exp\left(-\delta \sqrt{(-2b)(\phi_4 - 2(0))}\right)$$

$$\begin{aligned}
 &= C_3 \exp\left(-\delta\sqrt{(-2b)(\phi_4)}\right) \\
 &= -\sqrt{b}\sqrt{\frac{\pi}{2}} \exp(0)
 \end{aligned}$$

$\implies C_3 = -\sqrt{b}$ ,  $\sqrt{\phi_4(-2b)} = 0 \implies \phi_4 = 0$ , since  $b \neq 0$ .  $b \neq 0$  since  $F(a, b) = C_3$  does not satisfy the PDE.

$$\implies F(a, b) = \frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2b(\phi_4 - 2a)}\right)$$

or

$$F(a, b) = -\sqrt{b}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2a(\phi_2 - 2b)}\right)$$

Determining which of the two satisfies the PDE, let

$$F(a, b) = \frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2b(\phi_4 - 2a)}\right)$$

$$\begin{aligned}
 \implies \frac{\partial^2}{\partial a \partial b} &\left(\frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2b(\phi_4 - 2a)}\right)\right) \\
 &= \frac{1}{2}\sqrt{\frac{\pi}{2}} \left(\frac{\delta \exp\left(\delta(-\sqrt{4ab - 2b\phi_4})\right) (4a\delta\sqrt{b(2a - \phi_4)} - \sqrt{2}\phi_4)}{4a\sqrt{a}\sqrt{b(2a - \phi_4)}}\right).
 \end{aligned}$$

At  $\phi_4 = 0$ ,

$$\begin{aligned}
 \frac{\partial^2}{\partial a \partial b} &\left(\frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\delta\sqrt{ab}\right)\right) = \frac{1}{2}\sqrt{\frac{\pi}{2}} \left(\frac{\delta \exp\left(-\delta\sqrt{4ab}\right) (4a\delta)}{4a\sqrt{a}}\right) \\
 &= \frac{1}{2}\sqrt{\frac{\pi}{2}} \left(\frac{\delta^2 \exp\left(-\delta\sqrt{4ab}\right)}{\sqrt{a}}\right) = \frac{\delta^2}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\delta\sqrt{ab}\right)
 \end{aligned}$$

If  $\delta = 1$ , then  $\frac{\partial^2}{\partial a \partial b} \left(\frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\delta\sqrt{ab}\right)\right) = \frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\sqrt{ab}\right)$ .

Also,  $F(a, b) = \frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\sqrt{ab}\right)$ .

It implies that  $F(a, b) = \frac{1}{2\sqrt{a}}\sqrt{\frac{\pi}{2}} \exp\left(-2\sqrt{ab}\right)$  satisfies the PDE.

Also let  $F(a, b) = -\sqrt{b}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2a(\phi_2 - 2b)}\right)$

$$\begin{aligned}
 \implies \frac{\partial^2}{\partial a \partial b} &\left(-\sqrt{b}\sqrt{\frac{\pi}{2}} \exp\left(-\delta\sqrt{-2a(\phi_2 - 2b)}\right)\right) \\
 &= \sqrt{\frac{\pi}{2}} \left(\frac{\delta \exp\left(-\delta\sqrt{-2a(\phi_2 - 2b)}\right) (\sqrt{2}(4b - \phi_2) - 4b\delta\sqrt{a(2b - \phi_2)})}{4\sqrt{b}\sqrt{a(2b - \phi_2)}}\right).
 \end{aligned}$$

At  $\phi_2 = 0$ ,

$$\frac{\partial^2}{\partial a \partial b} \left(-\sqrt{b}\sqrt{\frac{\pi}{2}} \exp\left(-2\delta\sqrt{ab}\right)\right) = \left(-\delta^2\sqrt{b}\exp\left(-2\delta\sqrt{ab}\right) + \frac{\delta}{\sqrt{a}} \exp\left(-2\delta\sqrt{ab}\right)\right).$$

But

$$\sqrt{\frac{\pi}{2}} \left( -\delta^2 \sqrt{b} \exp(-2\delta\sqrt{ab}) + \frac{\delta}{\sqrt{a}} \exp(-2\delta\sqrt{ab}) \right) \neq -\sqrt{b} \sqrt{\frac{\pi}{2}} \exp(-2\delta\sqrt{ab})$$

for constant values of  $\delta$  and this implies that  $F(a, b) \neq -\sqrt{b} \sqrt{\frac{\pi}{2}} \exp(-2\delta\sqrt{ab})$  does not satisfy the PDE for all constant values of  $\delta$ .

Hence,

$$F(a, b) = \frac{1}{2\sqrt{a}} \sqrt{\frac{\pi}{2}} \exp(-2\sqrt{ab})$$

By the result in (11),

$$K_{\frac{1}{2}}(x) = \sqrt{\frac{\pi}{2x}} e^{-x}$$

$$\begin{aligned} \Rightarrow F(a, b) &= \frac{1}{2\sqrt{a}} \sqrt{\frac{\pi}{2}} \exp(-2\sqrt{ab}) = \sqrt{\sqrt{\frac{b}{2a}} \sqrt{\frac{\pi}{2 \cdot 2\sqrt{ab}}}} e^{-2\sqrt{ab}} \\ &= \left(\frac{b}{2a}\right)^{\frac{1}{4}} K_{\frac{1}{2}}\{2\sqrt{ab}\} \end{aligned}$$

$$\therefore \int_0^{\infty} \sin\left(au^2 - \frac{b}{u^2}\right) du = \left(\frac{b}{2a}\right)^{\frac{1}{4}} K_{\frac{1}{2}}\{2\sqrt{ab}\}$$

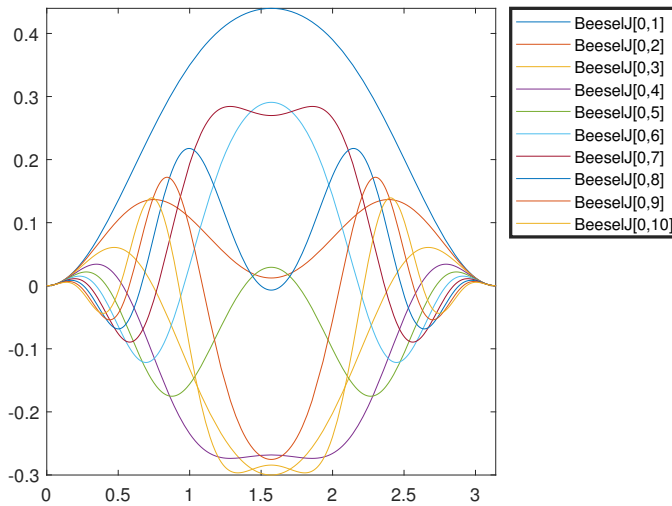
■

## 10 Main Results/Propositions

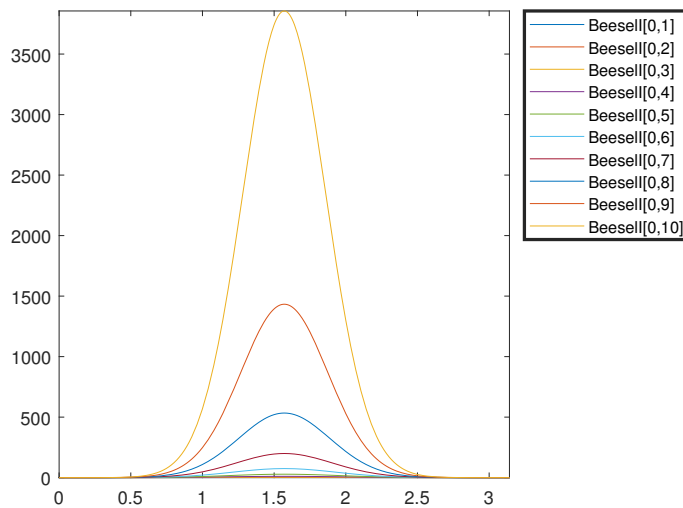
$$\begin{aligned}
 J_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 J_0(x\sqrt{1-t}) t^{\alpha-1} dt \\
 &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi 2 \sin^{2\alpha-1} \vartheta \sin \phi \frac{\partial}{\partial (x \cos \vartheta)} (x \cos \vartheta \sin(x \cos \vartheta \sin \phi)) d\phi d\vartheta \\
 I_\alpha(x) &= \frac{1}{\Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^1 I_0(x\sqrt{1-t}) t^{\alpha-1} dt \\
 &= \frac{1}{\pi x \Gamma(\alpha)} \left(\frac{x}{2}\right)^\alpha \int_0^{\frac{\pi}{2}} \int_0^\pi 2 \sin^{2\alpha-1} \vartheta \sin \phi \frac{\partial}{\partial (x \cos \vartheta)} (x \cos \vartheta \sinh(x \cos \vartheta \sin \phi)) d\phi d\vartheta \\
 J_0(x) &= \frac{1}{\pi x} \int_0^\pi \sin \phi \frac{\partial}{\partial x} (x \sin(x \sin \phi)) d\phi \\
 I_0(x) &= \frac{1}{\pi x} \int_0^\pi \sin \phi \frac{\partial}{\partial x} (x \sinh(x \sin \phi)) d\phi \\
 \int_0^\infty \sin\left(au^2 - \frac{b}{u^2}\right) du &= \left(\frac{b}{2a}\right)^{\frac{1}{4}} K_{\frac{1}{2}}\{2\sqrt{ab}\}
 \end{aligned}$$



**10.1 The area between each of the curves and the  $x$ -axis gives the values of  $J_0(z)$  from  $z = 1$  to 10**



**10.2 The area between each of the curves and the  $x$ -axis gives the values of  $I_0(z)$  from  $z = 1$  to 10**



## References

- [1] George Neville Watson Sc.D, F.R.S. *A TREATISE ON THE THEORY OF BESSEL FUNCTIONS*, chapter I, page 1. Cambridge University Press, 1944.
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- [4] George Neville Watson Sc.D, F.R.S. *A TREATISE ON THE THEORY OF BESSEL FUNCTIONS*, chapter VI, page 180. Cambridge University Press, 1944.
- [5] [https://en.wikipedia.org/wiki/Bessel\\_function](https://en.wikipedia.org/wiki/Bessel_function).